



**Daffodil**  
*International*  
**University**

# Microwave Engineering

## ETE 415

**LECTURE 3**  
**TRANSMISSION-LINE THEORY-II**

# Voltage Standing Wave Ratio

(lossless TL)

The total phasor voltage as a function of position on a line connected to a load at  $z=0$  is

$$V(-\ell) = V_o^+ e^{j\beta\ell} (1 + \Gamma_L e^{-j2\beta\ell})$$

Let  $\Gamma_L = |\Gamma_L| e^{j\theta}$

$$V(-\ell) = V_o^+ e^{j\beta\ell} (1 + |\Gamma_L| e^{-j(2\beta\ell - \theta)}) \quad \& \quad V^*(-\ell) = V_o^{+*} e^{-j\beta\ell} (1 + |\Gamma_L| e^{+j(2\beta\ell - \theta)})$$

$$|V(-\ell)| = \sqrt{V(-\ell)V^*(-\ell)}$$

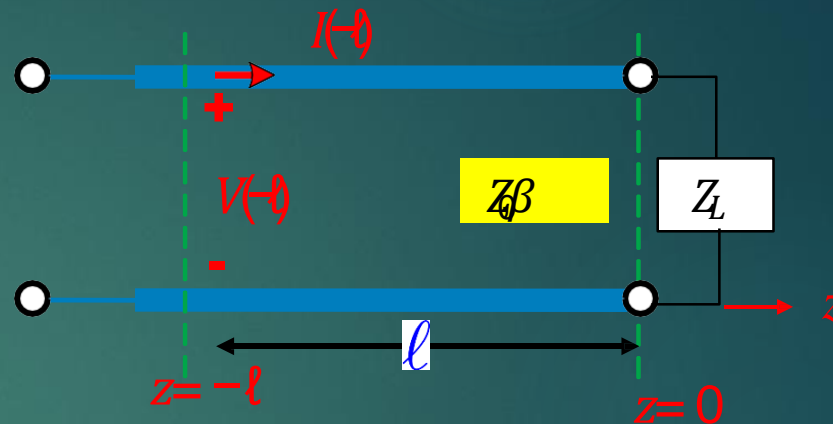
The amplitude of voltage as a function of  $z=-\ell$

$$= |V_o^+| \sqrt{(1 + |\Gamma_L| e^{-j(2\beta\ell - \theta)}) (1 + |\Gamma_L| e^{+j(2\beta\ell - \theta)})}$$

$$= |V_o^+| \sqrt{1 + |\Gamma_L| e^{+j(2\beta\ell - \theta)} + |\Gamma_L| e^{-j(2\beta\ell - \theta)} + |\Gamma_L|^2}$$

$$= |V_o^+| \sqrt{1 + 2\Re\{|\Gamma_L| e^{+j(2\beta\ell - \theta)}\} + |\Gamma_L|^2}$$

$$= |V_o^+| \sqrt{1 + |\Gamma_L|^2 + 2|\Gamma_L| \cos(2\beta\ell - \theta)}$$



# Voltage Standing Wave Ratio

(lossless TL)

$$|V|_{\max} = |V_o^+| \sqrt{1 + |\Gamma_L|^2 + 2|\Gamma_L|} = |V_o^+|(1 + |\Gamma_L|)$$


$$2\beta l_{\max,n} - \theta = 2n\pi$$

**$n$  is an integer**

$$|V|_{\min} = |V_o^+| \sqrt{1 + |\Gamma_L|^2 - 2|\Gamma_L|} = |V_o^+|(1 - |\Gamma_L|)$$

$$2\beta l_{\min,n} - \theta = (2n + 1)\pi$$

■ The voltage standing wave ratio (VSWR) is defined as

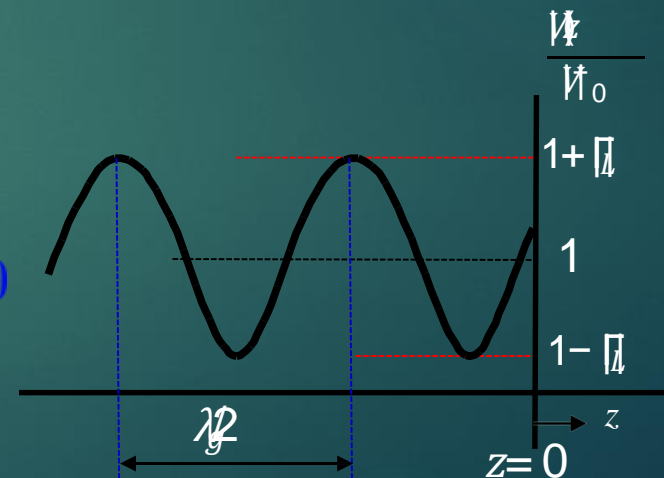
$$\text{VSWR} = \frac{V_{\max}}{V_{\min}} = \frac{1 + |\Gamma_L|}{1 - |\Gamma_L|} \geq 1$$


■ For two adjacent maxima at, say,  $n=1$  and  $n=0$  we can write

$$2\beta l_{\max,n=1} - \theta = 2\pi \text{ and } 2\beta l_{\max,n=0} - \theta = 0$$

$$\Delta l = l_{\max,n=1} - l_{\max,n=0} = \frac{\pi}{\beta} = \frac{\pi}{\frac{2\pi}{\lambda_g}} = \frac{\lambda_g}{2}$$

■ Voltage maxima and minima repeat every  $\lambda_g/2$



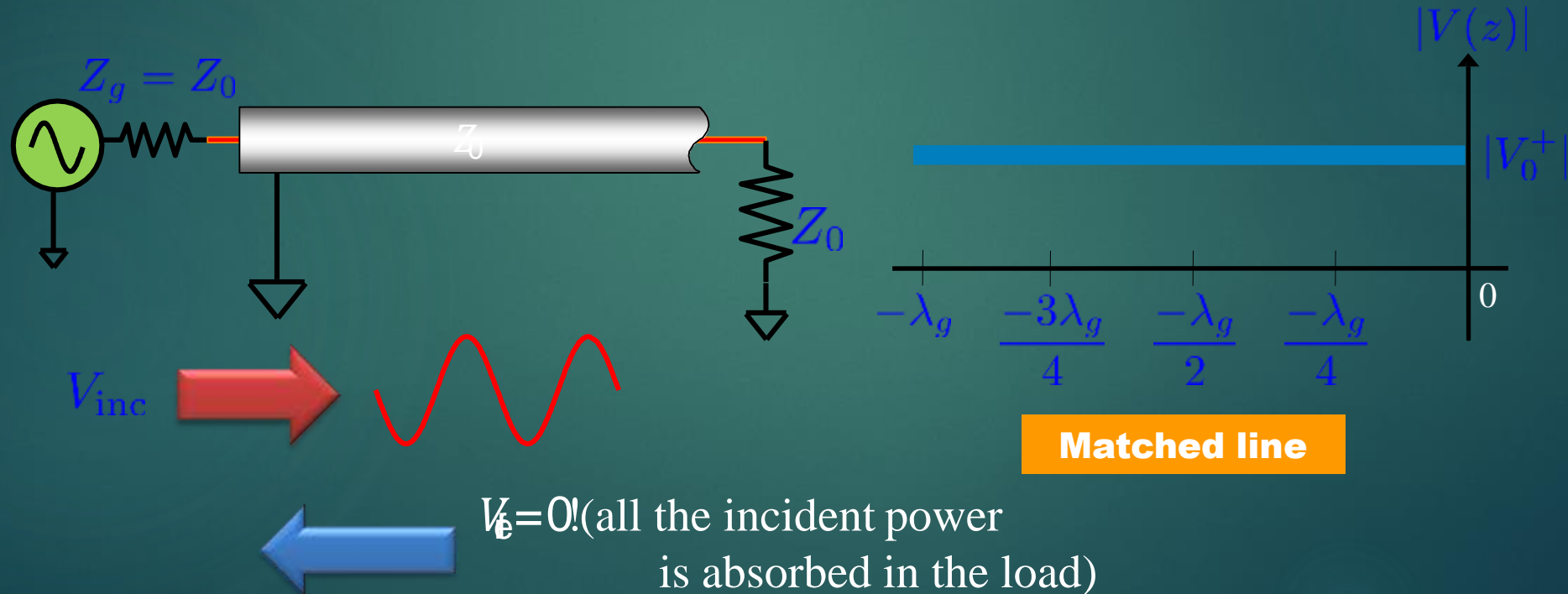
# Special cases

1. If  $Z_L = Z_0$  (matched load) then  $\Gamma_L = 0$ . Consequently,  $|\Gamma_L| = 0 \Rightarrow \text{VSWR} = 1$ .

$$V(z) = V_o^+ e^{-j\beta z} \Rightarrow |V(z)| = |V_o^+|$$

$$I(z) = \frac{V_o^+}{Z_0} e^{-j\beta z} \Rightarrow |I(z)| = \frac{|V_o^+|}{Z_0}$$

for all values of z

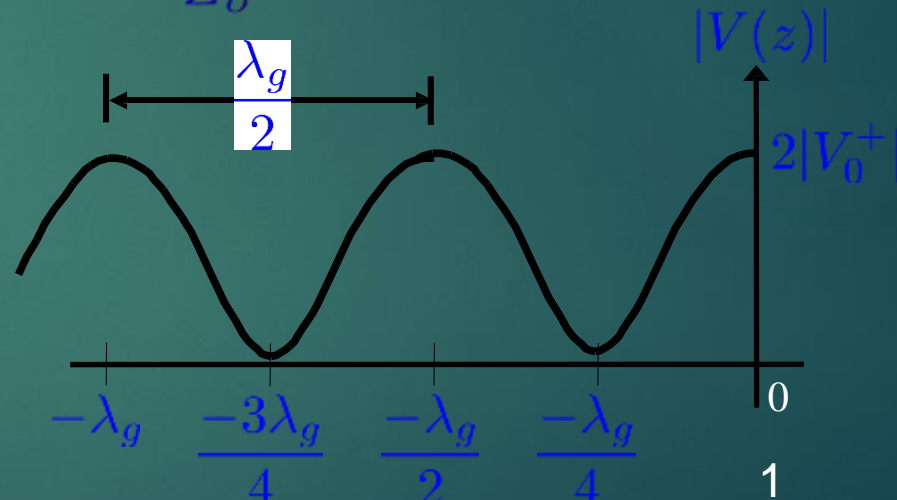
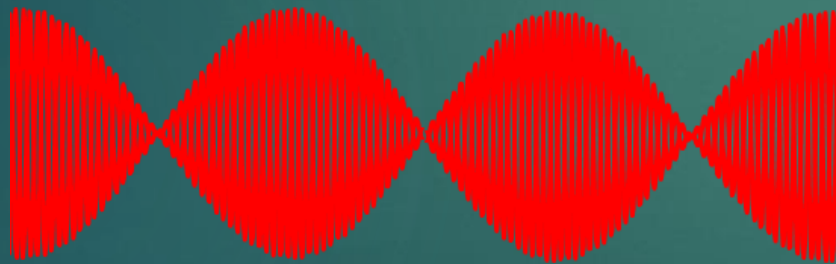


# Special cases

2. If  $Z_L = \infty$  (open-circuit load) then  $\Gamma_L = 1$ . Consequently,  
 $|\Gamma_L| = 1 \Rightarrow \text{VSWR} = \infty$ ,

$$V(z) = V_o^+ \left( e^{-j\beta z} + \underbrace{\frac{V_o^-}{V_o^+}}_{\equiv \Gamma_L = 1} e^{j\beta z} \right) = 2V_o^+ \cos(\beta z)$$

$$I(z) = \frac{V_o^+}{Z_0} \left( e^{-j\beta z} - \underbrace{\frac{V_o^-}{V_o^+}}_{\equiv \Gamma_L = 1} e^{j\beta z} \right) = -\frac{j2V_o^+}{Z_0} \sin(\beta z)$$



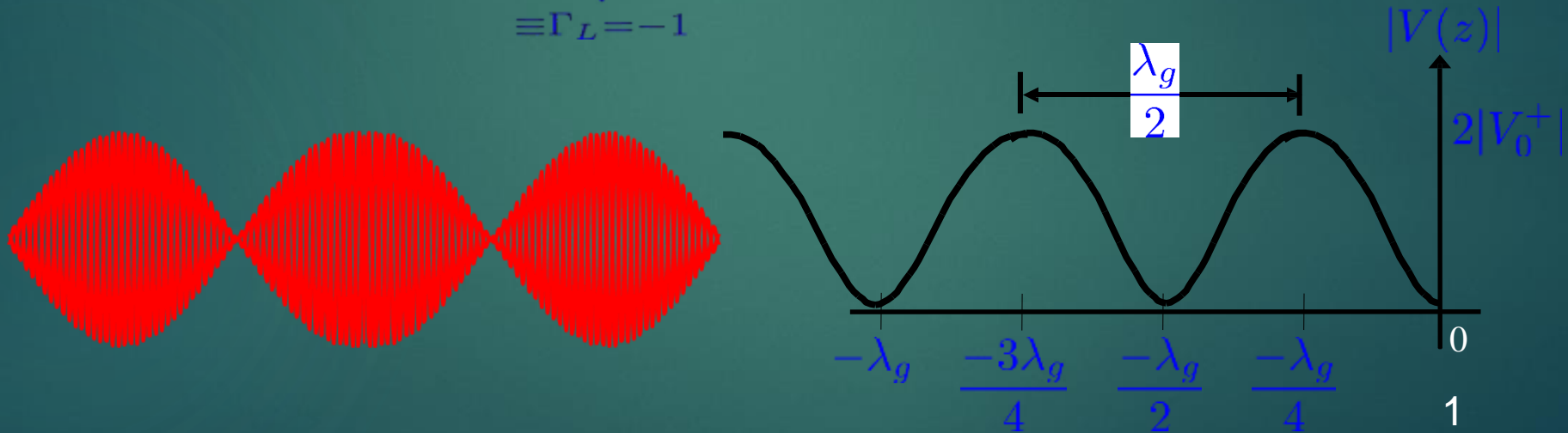
Open circuit

# Special cases

3. If  $Z_L = 0$  (short-circuit load) then  $\Gamma_L = -1$ . Consequently,  
 $|\Gamma_L| = 1 \Rightarrow \text{VSWR} = \infty$ ,

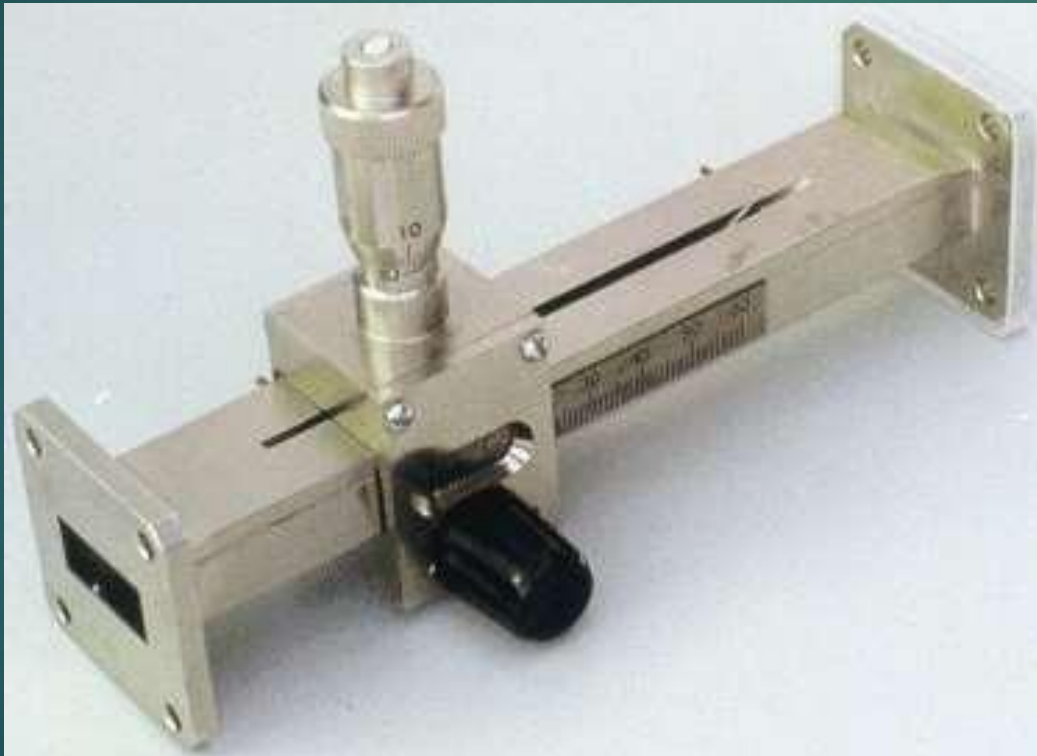
$$V(z) = V_o^+ \left( e^{-j\beta z} + \underbrace{\frac{V_o^-}{V_o^+}}_{\equiv \Gamma_L = -1} e^{j\beta z} \right) = -j2V_o^+ \sin(\beta z)$$

$$I(z) = \frac{V_o^+}{Z_0} \left( e^{-j\beta z} - \underbrace{\frac{V_o^-}{V_o^+}}_{\equiv \Gamma_L = -1} e^{j\beta z} \right) = \frac{2V_o^+}{Z_0} \cos(\beta z)$$



# The Slotted Line

■ A **slotted line** is a TL configuration (usually a waveguide or coaxial line) that allows the sampling of the electric field amplitude of a standing wave on a terminated line.



■ Used to measure:

- SWR
- $\lambda$
- $l_{\min}$   $l_{\max}$
- $Z_L$

# Slotted Line Measurement Procedure

- Calculate reflection coefficient magnitude from SWR

$$\text{VSWR} = \frac{V_{\max}}{V_{\min}} = \frac{1 + |\Gamma_L|}{1 - |\Gamma_L|} \Rightarrow |\Gamma_L| = \frac{\text{VSWR} - 1}{\text{VSWR} + 1}$$

- locations of voltage minima respective to load terminals are:

$$2\beta l_{\min,n} - \theta = (2n + 1)\pi$$

$$l_{\min,n} = (2n + 1) \frac{\lambda_g}{4} + \frac{\theta}{\pi} \cdot \frac{\lambda_g}{4}, n = 0, 1, \dots$$

$$\Rightarrow \theta = 2\beta l_{\min,n} - (2n + 1)\pi, n = 0, 1, \dots$$

$$2\beta = \frac{4\pi}{\lambda_g}$$

- For  $0 \leq l_{\min} \leq \frac{\lambda_g}{2}$  (1<sup>st</sup> minimum) and  $-\pi \leq \theta < \pi$

$$\theta = 2\beta l_{\min} - \pi$$

**See example 2.4 pp.**

**70-72**

- load impedance is calculated as

$$Z_L = Z_0 \frac{1 + \Gamma_L}{1 - \Gamma_L} = Z_0 \frac{1 + |\Gamma_L|e^{j\theta}}{1 - |\Gamma_L|e^{j\theta}}$$



# Time-Average Power Flow on TLs

■ A hugely important part of MW engineering is **delivering signal power to a load**.

■ **Time average power  $(P_{in})_{av}$**  delivered to input of TL, i.e.,  $z = -\ell$

$$(P_{in})_{av} = \frac{1}{2} \Re [V(-\ell) I^*(-\ell)]$$

similar to that used in circuit analysis.

$$V(-\ell) = V_0^+ (e^{\alpha\ell} e^{j\beta\ell} + \Gamma_L e^{-\alpha\ell} e^{-j\beta\ell}) \quad \& \quad I(-\ell) = \frac{V_0^+}{Z_0} (e^{\alpha\ell} e^{j\beta\ell} - \Gamma_L e^{-\alpha\ell} e^{-j\beta\ell})$$

$$I^*(-\ell) = \frac{V_0^{+*}}{Z_0} (e^{\alpha\ell} e^{-j\beta\ell} - \Gamma_L^* e^{-\alpha\ell} e^{+j\beta\ell})$$

Assume  $Z_0$  is real

$$(P_{in})_{av} = \frac{1}{2} \frac{|V_0^+|^2}{Z_0} \Re \left[ e^{2\alpha\ell} - |\Gamma_L| e^{+j(2\beta\ell - \theta)} + |\Gamma_L| e^{-j(2\beta\ell - \theta)} - |\Gamma_L|^2 e^{-2\alpha\ell} \right]$$

$$= \frac{1}{2} \frac{|V_0^+|^2}{Z_0} \Re \left[ e^{2\alpha\ell} - j2|\Gamma_L| \sin(\beta\ell - \theta) - |\Gamma_L|^2 e^{-2\alpha\ell} \right]$$

$$= (P_i)_{av} \left[ \underbrace{e^{2\alpha\ell}}_{\text{Forward}} - \underbrace{|\Gamma_L|^2 e^{-2\alpha\ell}}_{\text{Backward}} \right] \quad [\text{W}]$$

Forward

Backward

$$(P_i)_{av} = \frac{1}{2} \frac{|V_0^+|^2}{Z_0}$$

# Time-Average Power Flow on TLs

- Power delivered to load may be found by setting  $l = 0$

$$(P_L)_{av} = (P_i)_{av} (1 - |\Gamma_L|^2) \quad [\text{W}]$$



- For **entirely reactive** load,  $|\Gamma_L| = 1$ , then  $(P_L)_{av} = 0$  and no time average power is delivered to the load, as expected.

For all other passive loads,  $(P_L)_{av} > 0$ .

- the power dissipated in TL is the difference between the input power and the power delivered to the load

$$\begin{aligned} P_{\text{loss}} &= (P_{\text{in}})_{av} - (P_L)_{av} \\ &= (P_i)_{av} [|\Gamma_L|^2(1 - e^{-2\alpha l}) + e^{2\alpha l} - 1] \end{aligned}$$

- For lossless TL, i.e.,  $\alpha = 0$

$$(P_{\text{in}})_{av} = (P_L)_{av} = (P_i)_{av} (1 - |\Gamma_L|^2) \quad \Rightarrow \quad P_{\text{loss}} = 0$$

# Time-Average Power Flow on TLs

## Return loss (RL)

- The **relative reflected time average power** from an arbitrary load on a lossless TL  $= |\Gamma_L|^2$ .
- The relative time average power that is not delivered to the load can be considered a **“loss”**
  - ➔ since the signal from the generator was intended to be completely transported – not returned to the generator.

- This **return loss (RL)** is defined as

$$RL = -10 \log_{10} (|\Gamma_L|^2) = -20 \log_{10} (|\Gamma_L|) \quad \text{dB}$$

- The two extremes for return loss with a passive load are:
  1. **A matched load** where  $\Gamma_L = 0$  and  $RL = \infty$  dB (no reflected power)
  2. **A reactive load** where  $|\Gamma_L| = 1$  and  $RL = 0$  dB (all power reflected).

# Dispersion

(lossy TL)

From before, we know  $\beta = \frac{2\pi}{\lambda_g} = \frac{\omega}{v_p} \Rightarrow$

$$v_p = \frac{\omega}{\beta}$$

$$\gamma = \alpha + j\beta$$

For a lossy line,  $\beta = \text{Im}\{\sqrt{(R' + j\omega L')(G' + j\omega C')}\}$

Then,  $v_p = \frac{\omega}{\text{Im}\{\sqrt{(R' + j\omega L')(G' + j\omega C')}\}}$

**$v_p$  is dependent on frequency!!**

Is this a problem?



Signals that carry any information must occupy some bandwidth.

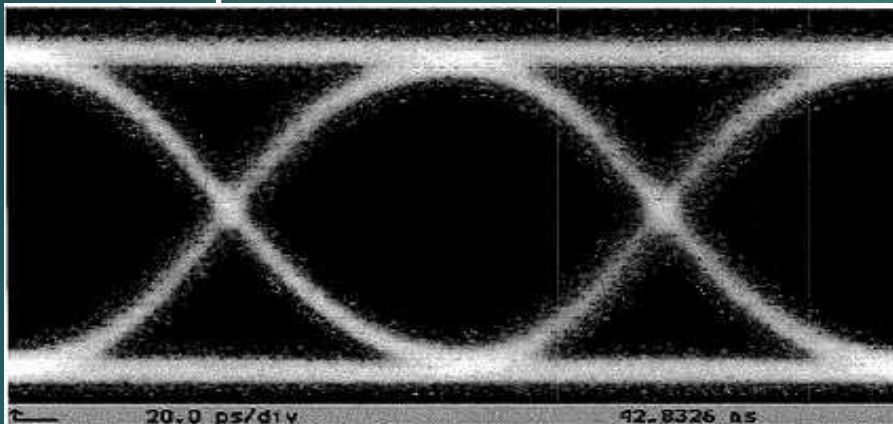
- ➔ FM broadcast (stereo)  $\Rightarrow$  occupies bandwidth  $\approx 100$  kHz
- ➔ AM analog modulation  $\Rightarrow$  occupies bandwidth  $\approx 6$  kHz

What about digital signals?

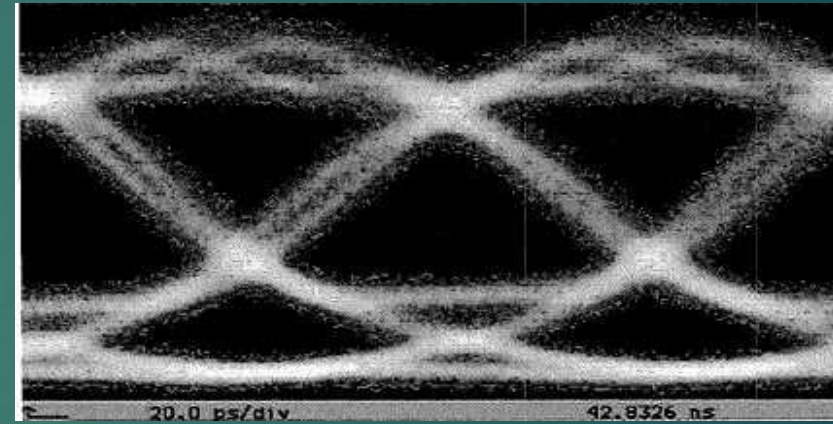
- ➔ A high-speed digital signal has  ~~$t_f$~~  on the order of 10's of ps.

# Dispersion

- Many high frequency components make up the edges.
- What if each frequency travels at a different velocity?
  - ▶ When each frequency making up an edge travels at a different speed, the waveform edges get ~~—smeared~~ over time and space.  
Dispersion!

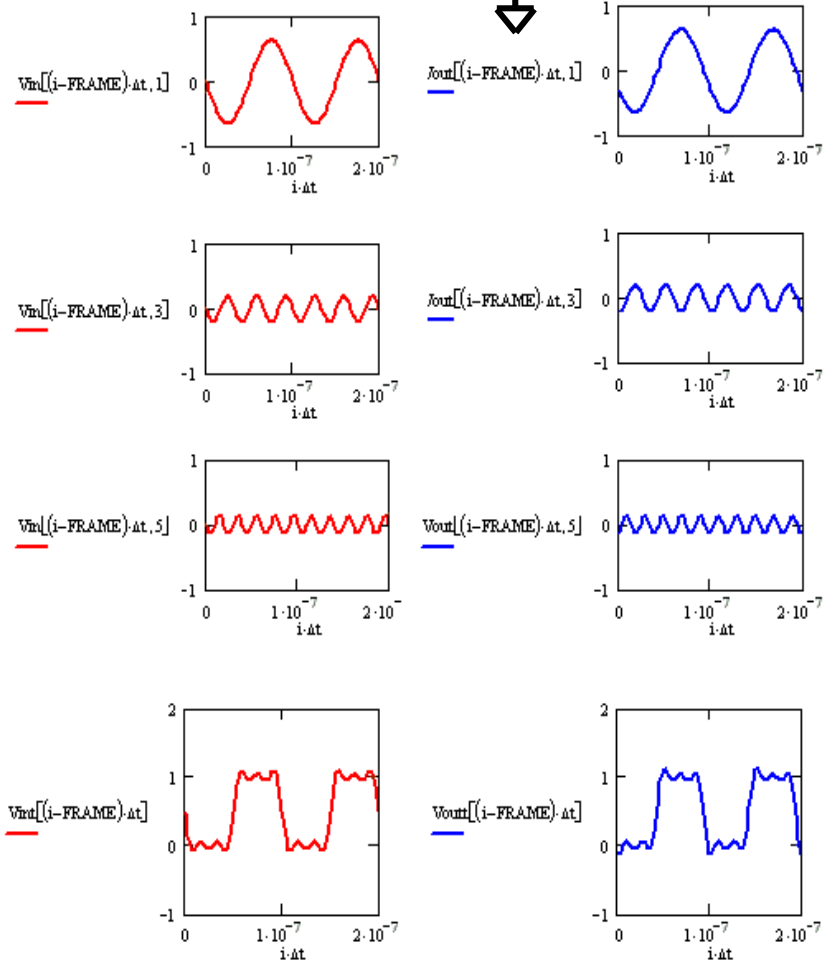


**Transmitted signal @  $\approx 5$  GHz “A very high-speed digital signal”**

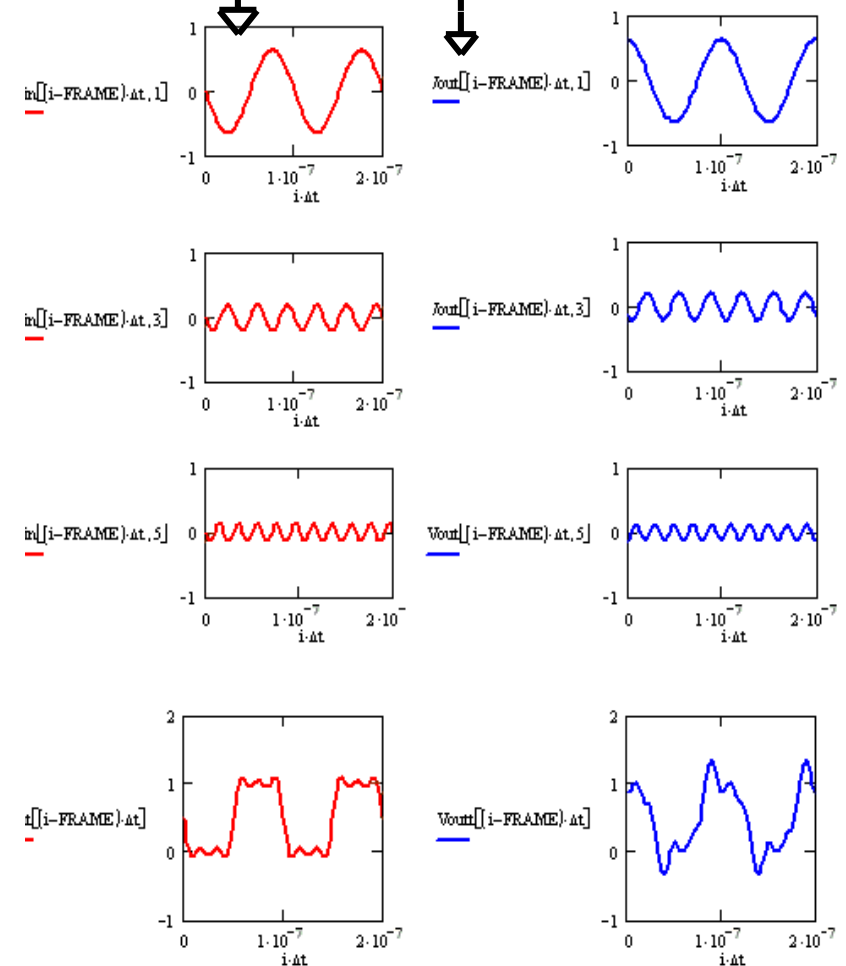


**Received signal after travelling down a lossy TL with loss.**

- Lossy lines can strongly affect high-speed digital signals.
- What about lower frequencies/longer lines?



**Low dispersion T-line**



**High dispersion T-line**

# Distortionless Line

■ The complex propagation constant

$$\gamma = \sqrt{(R' + j\omega L')(G' + j\omega C')}$$

■ Can be rearranged as

$$\begin{aligned}\gamma &= \sqrt{(j\omega L')(j\omega C') \left(1 + \frac{R'}{j\omega L'}\right) \left(1 + \frac{G'}{j\omega C'}\right)} \\ &= j\omega\sqrt{L'C'} \sqrt{1 - j\left(\frac{R'}{\omega L'} + \frac{G'}{\omega C'}\right) - \frac{R'G'}{\omega^2 L'C'}}\end{aligned}$$

■ For  $\frac{R'}{L'} = \frac{G'}{C'}$

$$\begin{aligned}\gamma &= j\omega\sqrt{L'C'} \sqrt{1 - 2j\frac{R'}{\omega L'} - \left(\frac{R'}{\omega L'}\right)^2} \\ &= j\omega\sqrt{L'C'} \left(1 - j\frac{R'}{\omega L'}\right) \\ &= R' \sqrt{\frac{C'}{L'}} + j\omega\sqrt{L'C'} = \alpha + j\beta\end{aligned}$$

$\beta = \omega \sqrt{L'C'} \Rightarrow$  **linear function of frequency**

$$v_p = \frac{\omega}{\omega\sqrt{L'C'}} = \frac{1}{\sqrt{L'C'}}$$

$\alpha = R' \sqrt{\frac{C'}{L'}} \Rightarrow$  **independent of frequency**

# The Low-Loss Line

- In most practical microwave and RF transmission lines the loss is small—if this were not the case, the line would be of little practical value.
- When the loss is small, some approximations can be made to simplify the expressions for the general transmission line parameters of  $\gamma = \alpha + j\beta$  and  $Z_0$

$$\gamma = j\omega\sqrt{L'C'}\sqrt{1 - j\left(\frac{R'}{\omega L'} + \frac{G'}{\omega C'}\right) - \frac{R'G'}{\omega^2 L'C'}}$$

- For a low-loss line both conductor and dielectric loss will be small, and we can assume that  $R' \ll \omega L'$  and  $G' \ll \omega C'$ . Then,  $R'G' \ll \omega^2 L'C'$ , and above equation reduces to

$$\gamma \simeq j\omega\sqrt{L'C'}\sqrt{1 - j\left(\frac{R'}{\omega L'} + \frac{G'}{\omega C'}\right)}$$
$$\alpha + j\beta \simeq j\omega\sqrt{L'C'}\left[1 - \frac{j}{2}\left(\frac{R'}{\omega L'} + \frac{G'}{\omega C'}\right)\right]$$

**Taylor approximation:**

$$\sqrt{1 \pm x} \approx 1 \pm \frac{x}{2}, \quad |x| \ll 1$$



# The Low-Loss Line

$$\alpha = \operatorname{Re}\{\gamma\}$$

$$\begin{aligned} &= \frac{w}{2} \sqrt{L'C'} \left( \frac{R'}{wL'} + \frac{G'}{wC'} \right) = \frac{1}{2} \left( \frac{\sqrt{L'C'} R'}{L'} + \frac{\sqrt{L'C'} G'}{C'} \right) \\ &= \frac{1}{2} \left( R' \sqrt{\frac{C'}{L'}} + G' \sqrt{\frac{L'}{C'}} \right) = \frac{1}{2} \left( \frac{R'}{Z_0} + G' Z_0 \right) \end{aligned}$$

Higher  $Z_0$  increases dielectric loss and reduces conduction loss

$$\beta = \operatorname{Im}\{\gamma\} = w \sqrt{L'C'} \quad \text{Same as no loss line} \Rightarrow \text{negligible dispersion}$$



the characteristic impedance  $Z_0$  can be approximated as a real quantity:

$$Z_0 = \sqrt{\frac{R' + jwL'}{G' + jwC'}} \simeq \sqrt{\frac{L'}{C'}}$$

Thank you Very Much !!!