

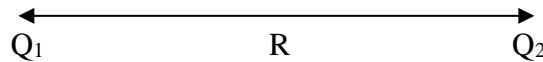
## CHAPTER-02

### Coulomb's Law and Electric Field Intensity

#### Coulomb's law

Coulomb's law states that the force  $F$  between two point charges  $Q_1$  and  $Q_2$  along the line joining them is-

- Directly proportional to the product  $Q_1Q_2$  of the charges
- Inversely proportional to the square of the distance  $R$  between



$$F \propto Q_1$$

$$F \propto Q_2$$

$$F \propto \frac{1}{r^2}$$

Now we can easily write an equation based on above equations.

$$F \propto \frac{Q_1 Q_2}{R^2} a_R$$

$$F = K \frac{Q_1 Q_2}{R^2} a_R \quad [\text{Where } K \text{ is proportional constant}]$$

Where,

$$\text{The value of, } K = \frac{1}{4\pi\epsilon_0}$$

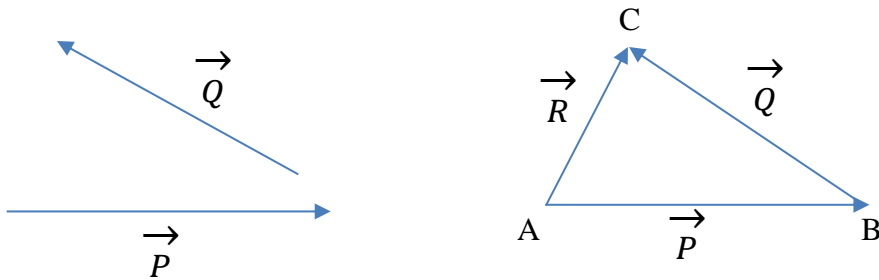
$$\therefore F = \frac{Q_1 Q_2}{4\pi\epsilon_0 R^2} a_R \quad [\text{Where, } \epsilon_0 = 8.854 \times 10^{-12} = \frac{1}{36\pi} \times 10^{-9} F/m]$$

The new constant  $\epsilon_0$  is called the permittivity of free space and has the magnitude, measured in farads per meter (F/m). This is a mathematical quantity that represents, how much electric field is permitted (Penetrated) in free space or vacuum. It is an ideal physical constant that represents the absolute dielectric permittivity of a vacuum. In other words,  $\epsilon_0$  quantifies the ability of a vacuum to allow electric field lines to flow through.

In SI units,  $Q_1$  and  $Q_2$  are measured in coulombs (C), the distance  $R$  is in meters (m) and the force  $F$  is in newtons (N).

## Triangle law of forces

If two vectors are represented in magnitude and direction by the two sides of a triangle taken in order, the third side of the triangle represents their resultant in magnitude and direction in reverse order.



- $\vec{P}$  and  $\vec{Q}$  are two non-zero forces
- These two are represented by two sides of a triangle AB, BC respectively
- The third side of the triangle AC is the resultant in the opposite direction.

**Mathematical problem-1:**

The use of the vector form of Coulomb's law consider a charge of  $Q_1 = 3 \times 10^{-4} C$  at P (1, 2, 3) and a charge of  $Q_2 = -10^{-4} C$  at Q (2, 0, 5) in a vacuum. Evaluate the force exerted on  $Q_2$  by  $Q_1$ .

**Solution:**

Given,

$$Q_1 = 3 \times 10^{-4} C$$

$$Q_2 = -10^{-4} C$$

Now,

$$\begin{aligned} R_{12} &= r_2 - r_1 \\ &= (2 - 1)a_x + (0 - 2)a_y + (5 - 3)a_z \\ &= a_x - 2 a_y + 2 a_z \end{aligned}$$

$$|R_{12}| = \sqrt{1^2 + (-2)^2 + 2^2} = \sqrt{9} = 3$$

$$a_{12} = \frac{1}{3} (a_x - 2 a_y + 2 a_z)$$

Thus,

$$\begin{aligned} F_2 &= \frac{Q_1 Q_2}{R^2} a_{12} \\ &= \frac{3 \times 10^{-4} \times (-10^{-4})}{4\pi \left(\frac{1}{36\pi} \times 10^{-9}\right) \times 3^2} \left(\frac{a_x - 2 a_y + 2 a_z}{3}\right) \\ &= -10(a_x - 2 a_y + 2 a_z) \\ &= -10a_x + 20a_y - 20a_z \quad (\text{Answer}) \end{aligned}$$

The force expressed by Coulomb's law is a mutual force, for each of the two charges experiences a force of the same magnitude, although of opposite direction.

$$F_1 = -F_2 = \frac{Q_1 Q_2}{4\pi\epsilon_0 (R_{12})^2} a_{21} = - \frac{Q_1 Q_2}{4\pi\epsilon_0 (R_{12})^2} a_{12}$$

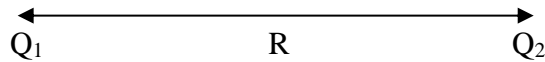
## Electric Field Intensity

The electric field intensity (or electric field strength)  $E$  is the force per unit charge when placed in the electric field. Thus

$$E = \lim_{Q \rightarrow 0} \frac{F}{Q}$$

Or simply

$$E = \frac{F}{Q}$$



Now consider one charge fixed in position, say  $Q_1$  and move a second charge slowly around, note that there exists everywhere a force on this second charge; in other words, this second charge is displaying the existence of a force field. Call this second charge a test charge  $Q_t$ . The force on it is given by Coulomb's law,

$$F_t = \frac{Q_1 Q_t}{R^2} a_{1t}$$

This force as a force per unit charge gives

$$\frac{F_t}{Q_t} = \frac{Q_1}{4\pi\epsilon_0 R_{1t}^2} a_{1t}$$

The quantity on the right side of the above equation is function only of  $Q_1$  and the directed line segment from  $Q_1$  to the position of the test charge. This describes a vector field and is called the electric field intensity.

We define the electric field intensity as the vector force on a unit positive test charge. Electric field intensity must be measured by the unit newtons per coulomb-the force per unit charge. We shall at once measure electric field intensity in the practical units of volts per meter (V/m). Using a capital letter  $\mathbf{E}$  for electric field intensity, we have

$$E = \frac{F_t}{Q_t}$$
$$E = \frac{Q_1}{4\pi\epsilon_0 R_{1t}^2} a_{1t}$$

Finally,

$$E = \frac{Q}{4\pi\epsilon_0 R^2} a_R$$

**Mathematical problem-2:**

Point charges 1 mC and -2 mC are located at (3, 2, -1) and (-1, -1, 4) respectively. Calculate the electric force on a 10 nC charge located at (0, 3, 1) and the electric field intensity at that point.

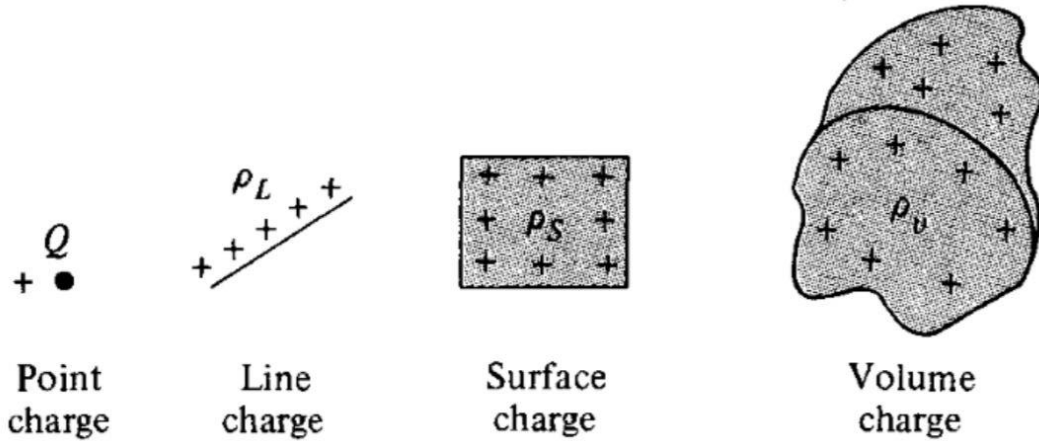
**Solution:**

$$\begin{aligned}
 F &= \sum_{K=1,2} \frac{Q Q_k}{4\pi\epsilon_0 R^2} a_R \\
 &= \sum_{K=1,2} \frac{Q Q_k (r - r_k)}{4\pi\epsilon_0 |r - r_k|^3} \\
 &= \frac{Q Q_1 (r - r_1)}{4\pi\epsilon_0 |r - r_1|^3} + \frac{Q Q_2 (r - r_2)}{4\pi\epsilon_0 |r - r_2|^3} \\
 &= \frac{Q}{4\pi\epsilon_0} \left\{ \frac{Q_1 (r - r_1)}{|r - r_1|^3} + \frac{Q_2 (r - r_2)}{|r - r_2|^3} \right\} \\
 &= \frac{Q}{4\pi\epsilon_0} \left\{ \frac{10^{-3} [(0, 3, 1) - (3, 2, -1)]}{|(0, 3, 1) - (3, 2, -1)|^3} - \frac{2 \cdot 10^{-3} [(0, 3, 1) - (-1, -1, 4)]}{|(0, 3, 1) - (-1, -1, 4)|^3} \right\} \\
 &= 9 \times 10^{-2} \left[ \frac{(-3, 1, 2)}{14\sqrt{14}} + \frac{(-2, -8, 6)}{26\sqrt{26}} \right] \\
 F &= -6.507 a_x - 3.817 a_y + 7.506 a_z \text{ mN}
 \end{aligned}$$

At that point,

$$\begin{aligned}
 E &= \frac{F}{Q} \\
 &= (-6.507, -3.817, 7.506) \times \frac{10^{-3}}{10 \times 10^{-9}} \\
 E &= -650.7 a_x - 381.7 a_y + 750.6 a_z \text{ kV/m}
 \end{aligned}$$

## Electric Fields due to continuous charge distributions



It is customary to denote the line charge density, surface charge density and volume charge density by  $\rho_L$  (in C/m),  $\rho_S$  (in C/m<sup>2</sup>) and  $\rho_V$  (in C/m<sup>3</sup>) respectively. The charge element  $dQ$  and the total charge  $Q$  due to these charge distributions are obtained from above figure as

$$dQ = \rho_L dl \rightarrow Q = \int_L \rho_L dl \text{ (Line charge)}$$

$$dQ = \rho_S dS \rightarrow Q = \int_S \rho_S dS \text{ (Surface charge)}$$

$$dQ = \rho_V dv \rightarrow Q = \int_V \rho_V dv \text{ (Volume charge)}$$

The electric field intensity due to each of the charge distributions  $\rho_L$ ,  $\rho_S$  and  $\rho_V$  may be regarded as the summation of the field contributed by the numerous point charges making  $\rho_L dl$ ,  $\rho_S dS$ ,  $\rho_V dv$  and integrating, we get

$$E = \int \frac{\rho_L dl}{4\pi\epsilon_0 R^2} a_R \text{ (Line charge)}$$

$$E = \int \frac{\rho_S dS}{4\pi\epsilon_0 R^2} a_R \text{ (Surface charge)}$$

$$E = \int \frac{\rho_V dv}{4\pi\epsilon_0 R^2} a_R \text{ (Volume charge)}$$

## Field due to a continuous volume charge distribution

### Cylindrical Co-ordinate

$$dQ = \rho_v dv$$

$$\int dQ = \int_v \rho_v dv$$

$$= \int_{z=0}^h \int_{\varphi=0}^{2\pi} \int_0^r \rho_v \rho d\rho d\varphi dz$$

### Spherical Co-ordinate

$$dQ = \rho_v dv$$

$$\int dQ = \int_v \rho_v dv$$

$$= \int_{\varphi=0}^{2\pi} \int_{\theta=0}^{\pi} \int_0^r \rho_v r^2 \sin\theta d\theta dr d\varphi$$

### Surface of Cylindrical Co-ordinate

$$\int dQ = \int_s \rho_s ds$$

$$Q = \int_0^{2\pi} \int_0^{\pi} \rho_s \rho d\rho d\varphi a_z$$

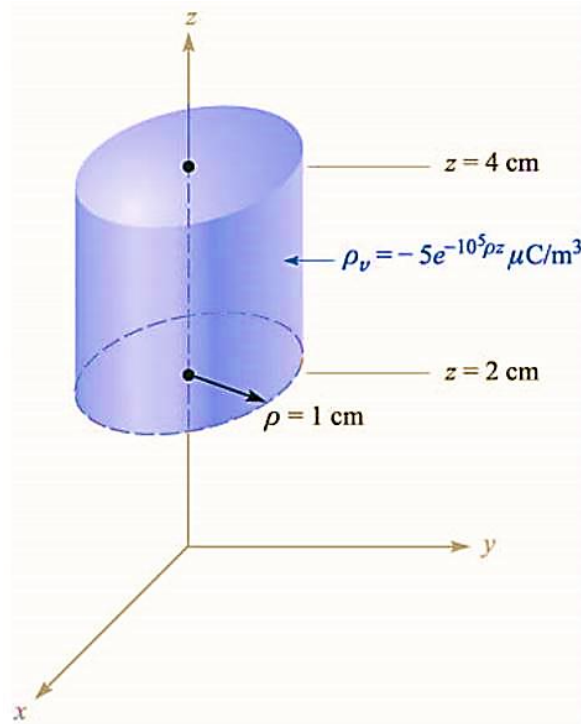
### Spherical Co-ordinate:

$$\int dQ = \int_s \rho_s ds$$

$$Q = \int_0^{2\pi} \int_0^{\pi} \rho_s (r^2 \sin\theta d\theta d\varphi) a_r$$

### Mathematical problem-3:

Evaluate the total charge contained in a 2-cm length of the electron beam show in following figure.



### Solution:

The total charge contained may be obtained by evaluating

$$Q = \int_v \rho_v dv$$

The charge density is

$$\rho_v = -5 \times 10^{-6} e^{-10^5} \text{ C}/\text{m}^2$$

Therefore,

$$\begin{aligned} Q &= \int_{0.02}^{0.04} \int_0^{2\pi} \int_0^{0.01} -5 \times 10^{-6} e^{-10^5 \rho z} \rho d\rho d\phi dz \\ &= -10^{-5} \pi \int_{0.02}^{0.04} \int_0^{0.01} e^{-10^5 \rho z} \rho d\rho dz \\ &= -10^{-5} \pi \int_0^{0.01} \left[ \frac{e^{-10^5 \rho z}}{10^{-5} \rho} \right]_{0.02}^{0.04} \rho d\rho \\ &= -10^{-10} \pi \int_0^{0.01} [e^{-2000\rho} - e^{-4000\rho}] d\rho \end{aligned}$$



$$= -10^{-10}\pi \left( \frac{1}{2000} - \frac{1}{4000} \right)$$

$$= -10^{-10}\pi \left( \frac{2-1}{4000} \right)$$

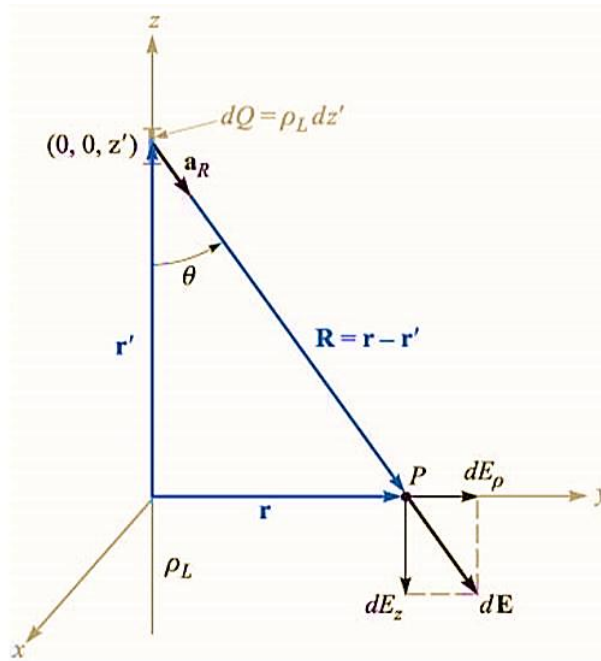
$$= -10^{-10}\pi \left( \frac{1}{40 \times 10^2} \right)$$

$$= \frac{-\pi}{40} \rho c$$

$$= 0.0785 \text{ pC} \quad (\text{Ans})$$

## Field of a Line Charge

Let us assume a straight line charge extending along the z-axis in a cylindrical co-ordinate system from  $-\infty$  to  $\infty$  as shown in following figure. We desire the electric field intensity,  $E$  at any and every point resulting from a uniform line charge density  $\rho_L$ .



As we move around the line charge, varying  $\phi$  while keeping  $\rho$  and  $z$  constant, the line charge appears the same from every angle. Again, if we maintain  $\rho$  and  $\phi$  constant while moving up and down the line charge by changing  $z$ , the line charge still recedes into infinite distance in both directions and the problem is unchanged. This is axial symmetry and leads to fields are not functions of  $z$ .

If we maintain  $\phi$  and  $z$  constant and vary  $\rho$ , the problem changes and Coulomb's law leads us to expect the field to become weaker as  $\rho$  increases. Hence, by a process of elimination we are led to the fact that the field varies only with  $\rho$ .

No element of charge produces a  $\phi$  component of electric intensity;  $E_\phi$  is zero. However, each element does produce an  $E_\rho$  and an  $E_z$  component but the contribution to  $E_z$  by elements of charge which are equal distances above and below the point at which we are determining the field will cancel.

Therefore we have found that we have only an  $E_\rho$  component and it varies only with  $\rho$ . Now to find this component we choose a point  $P(0, y, 0)$  on the  $y$ -axis at which to determine the field. This is perfectly general point in view of the lack of variation of the field with  $\phi$  and  $z$ . To find the incremental field at  $P$  due to the incremental charge  $dQ = \rho_L dz'$ , we have

$$dE = \frac{\rho_L dz' (r - r')}{4\pi\epsilon_0 |r - r'|^3}$$

Where

$$r = y\mathbf{a}_y = \rho\mathbf{a}_\rho$$

$$r' = z'\mathbf{a}_z$$

And

$$r - r' = \rho\mathbf{a}_\rho - z'\mathbf{a}_z$$

Therefore,

$$dE = \frac{\rho_L dz' (\rho\mathbf{a}_\rho - z'\mathbf{a}_z)}{4\pi\epsilon_0 (\rho^2 + z'^2)^{3/2}}$$

Since only the  $E_\rho$  component is present, we may simplify:

$$dE_\rho = \frac{\rho_L \rho dz'}{4\pi\epsilon_0 (\rho^2 + z'^2)^{3/2}}$$

And

$$E_\rho = \int_{-\infty}^{\infty} \frac{\rho_L \rho dz'}{4\pi\epsilon_0 (\rho^2 + z'^2)^{3/2}}$$

Integrating by integral tables or change of variable,  $z' = \rho \cot\theta$ , we have

$$E_\rho = \frac{\rho_L}{4\pi\epsilon_0} \rho \left( \frac{1}{\rho^2} \frac{z'}{\sqrt{\rho^2 + z'^2}} \right)_{-\infty}^{\infty}$$

$$E_\rho = \frac{\rho_L}{2\pi\epsilon_0 \rho}$$

$$E_\rho = \frac{\rho_L}{2\pi\epsilon_0 \rho} \mathbf{a}_\rho$$

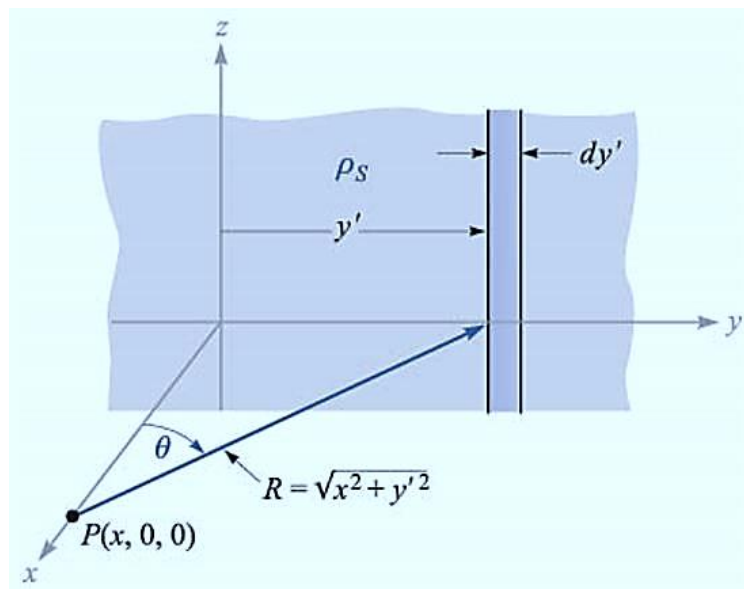
## Field of a Sheet of Charge

The infinite sheet of charge having a unit form density of  $\rho_s \text{ C/m}^2$ .

Such a charge distribution may often be used to approximate that found on the conductors of a strip transmission line or parallel-plate capacitor.  $\rho_s$  is commonly known as *surface charge density*.

Let us place a sheet of charge in the  $yz$  plane and again consider symmetry as following figure. The field cannot vary with  $y$  or with  $z$  and then that the  $y$  and  $z$  components arising from differential elements of charge symmetrically located with respect to the point at which we wish the field will cancel. Hence only  $E_x$  is present and this component is a function of  $x$  alone.

Let us use the field of the infinite line charge by dividing the infinite sheet into differential-width strips shown in following figure. The line charge density or charge per unit length is  $\rho_L = \rho_s dy'$  and the distance from this line charge to our general point P on the  $x$ -axis is  $R = \sqrt{x^2 + y'^2}$ .



$$E = \frac{\rho_L}{2\pi\epsilon_0 R} a_R$$

$$dE = \frac{\rho_s dy'}{2\pi\epsilon_0 R} a_R$$

The contribution to  $E_x$  at P from this differential-width strip is then

$$dE_x = \frac{\rho_s dy'}{2\pi\epsilon_0 \sqrt{x^2 + y'^2}} \cos\theta = \frac{\rho_s}{2\pi\epsilon_0} \frac{x dy'}{x^2 + y'^2}$$

Adding the effects of all the strips,

$$\begin{aligned} E_x &= \frac{\rho_s}{2\pi\epsilon_0} \int_{-\infty}^{\infty} \frac{x dy'}{x^2 + y'^2} \\ &= \frac{\rho_s}{2\pi\epsilon_0} \tan^{-1} \frac{y'}{x} \Big|_{-\infty}^{\infty} \end{aligned}$$

$$= \frac{\rho_s}{2\epsilon_0}$$

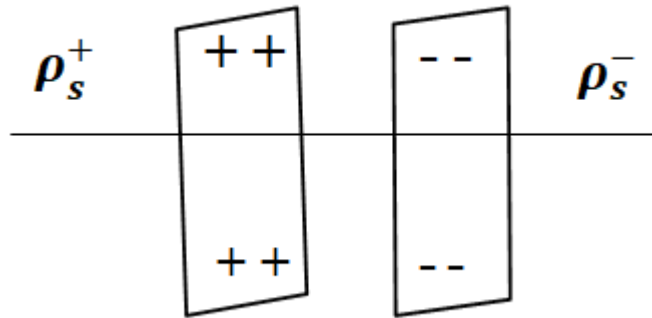
If the point P were chosen on the negative x-axis, then

$$E_x = -\frac{\rho_s}{2\epsilon_0}$$

For the field is always directed away from the positive charge. This difficulty in sign is usually overcome by specifying a unit vector  $a_N$ , which is normal to the sheet and directed outward or away from it. Then

$$E = \frac{\rho_s}{2\epsilon_0} a_N$$

If a second infinite sheet of charge having a negative charge density  $-\rho_s$ , is located in the plane  $x = a$ , we may find the total field by adding the contribution of each sheet.



In the region  $x > a$ ,

$$E_+ = \frac{\rho_s}{2\epsilon_0} a_x$$

$$E_- = -\frac{\rho_s}{2\epsilon_0} a_x$$

$$E = E_+ + E_- = 0$$

And for  $x < 0$ ,

$$E_+ = -\frac{\rho_s}{2\epsilon_0} a_x$$

$$E_- = \frac{\rho_s}{2\epsilon_0} a_x$$

$$E = E_+ + E_- = 0$$

And when  $0 < x < a$ ,

$$E_+ = \frac{\rho_s}{2\epsilon_0} a_x$$

$$E_- = \frac{\rho_s}{2\epsilon_0} a_x$$

$$E = E_+ + E_- = \frac{\rho_s}{\epsilon_0} a_x$$