

CHAPTER-4

Energy and potential

Energy expended in moving a Point charge in an Electric Field

We wish to move a charge Q a distance dL from (B) to (A) location in an electric field E . The force on Q due to the electric field is

$$F_E = QE$$

The force which we must apply is equal and opposite to the force due to the field. Therefore, the work done is,

$$F \cdot dL = -QE \cdot dL a_L$$

Where a_L = a unit vector in the direction of dL .

The differential work done by external source moving Q ,

$$= -QE \cdot dL a_L$$

$$= -QE \cdot dL$$

$$\therefore dW = -QE \cdot dL$$

- *If E and L are perpendicular, the differential work will be zero.*

The total work required to move the charge from (B) to (A) location is,

$$W = dW$$

$$\Rightarrow W = \int_B^A -QE \cdot dL$$

$$\therefore W = -Q \int_B^A E \cdot dL$$

- *$W > 0$ means we expend energy or do work.*
- *$W < 0$ means the field expends energy or do work.*

Mathematical problem-1:

The non-uniform field $E = y a_x + x a_y + 2 a_z$.

- (a) Determine the work expended in carrying 2 C from B(1, 0, 1) to A(0.8, 0.6, 1) along the shorter arc of the circle $x^2 + y^2 = 1$ and $z = 1$.
- (b) Determine the work required to carry 2 C from B to A in the same field, but this time use straight-line path from B to A.

Solution:

- (a) Working in Cartesian co-ordinates, the differential path dL is $dx a_x + dy a_y + dz a_z$ and the integral becomes

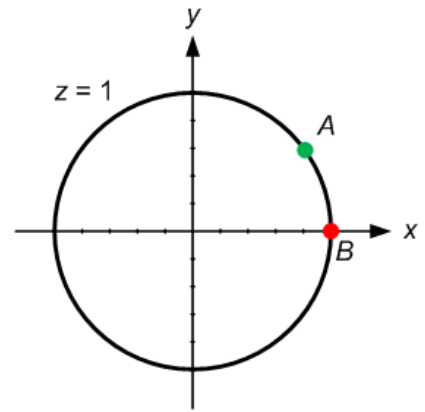
$$\begin{aligned} W &= -Q \int_B^A E \cdot dL \\ &= -2 \int_B^A (y a_x + x a_y + 2 a_z) \cdot (dx a_x + dy a_y + dz a_z) \\ &= -2 \int_1^{0.8} y dx - 2 \int_0^{0.6} x dy - 4 \int_1^1 dz \\ &= -2 \int_1^{0.8} \sqrt{1-x^2} dx - 2 \int_0^{0.6} \sqrt{1-y^2} dy - 0 \\ &\quad \text{[Circle equation: } x^2 + y^2 = 1, x = \sqrt{1-y^2}, y = \sqrt{1-x^2}] \\ &= -\left[x\sqrt{1-x^2} + \sin^{-1} x \right]_1^{0.8} - \left[y\sqrt{1-y^2} + \sin^{-1} y \right]_0^{0.6} \\ &\quad \left[\int \sqrt{a^2 - u^2} du = \frac{u}{2} \sqrt{a^2 - u^2} + \frac{a^2}{2} \sin^{-1} \frac{u}{a} \right] \\ &= -(0.48 + 0.927 - 0 - 1.571) - (0.48 + 0.644 - 0 - 0) \\ &= -0.96 J \quad (\text{Ans}) \end{aligned}$$

(b) The equations of the straight line,

$$y - y_B = \frac{y_A - y_B}{x_A - x_B} (x - x_B)$$

$$z - z_B = \frac{z_A - z_B}{y_A - y_B} (y - y_B)$$

$$x - x_B = \frac{x_A - x_B}{z_A - z_B} (z - z_B)$$



From the first equation above we have

$$y = -3x + 3$$

$$y = -3(x - 1)$$

And from the second we obtain

$$z = 1$$

Thus,

$$\begin{aligned} W &= -2 \int_1^{0.8} y \, dx - 2 \int_0^{0.6} x \, dy - 4 \int_1^1 dz \\ &= 6 \int_1^{0.8} (x - 1) \, dx - 2 \int_0^{0.6} \left(1 - \frac{y}{3}\right) dy \\ &= -0.96 \, J \quad (\text{Ans}) \end{aligned}$$

Differential Length

$$d\mathbf{L} = dx\mathbf{a}_x + dy\mathbf{a}_y + dz\mathbf{a}_z \quad (\text{Rectangular})$$

$$d\mathbf{L} = d\rho\mathbf{a}_\rho + \rho d\phi\mathbf{a}_\phi + dz\mathbf{a}_z \quad (\text{Cylindrical})$$

$$d\mathbf{L} = dr\mathbf{a}_r + r d\theta\mathbf{a}_\theta + r \sin \theta d\phi\mathbf{a}_\phi \quad (\text{Spherical})$$

Electric potential

We know, the total work required to move the charge from (B) to (A) location is,

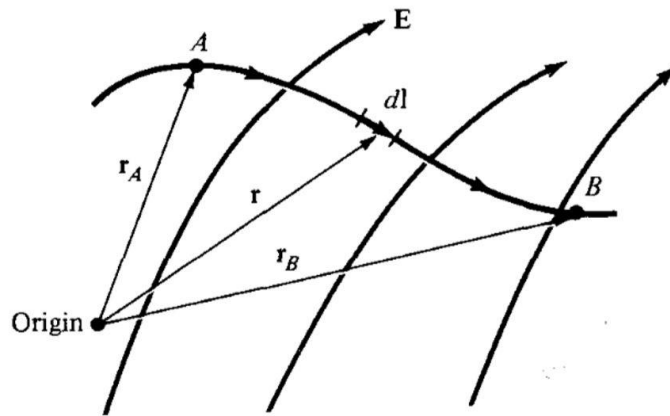
$$W = -Q \int_B^A E \cdot dL$$

Now dividing W by Q gives the potential energy per unit charge. The quantity denoted by V_{AB} is known as the potential difference between points B and A. Thus,

$$V_{AB} = \frac{W}{Q} = - \int_B^A E \cdot dL$$

Note that,

- In determining V_{AB} , B is the initial point while A is the final point.
- If V_{AB} is negative, there is a loss in potential energy in moving Q from B to A; this implies that the work is being done by the field. However, if V_{AB} is positive, there is a gain in potential energy in the movement.
- V_{AB} is measured in joules per coulomb, commonly referred to as volts (V).



If the E field in above figure is due to a point charge Q located at the origin, then

$$E = \frac{Q}{4\pi\epsilon_0 r^2} a_r$$

Then we have,

$$\begin{aligned} V_{AB} &= - \int_B^A E \cdot dL \\ V_{AB} &= - \int_{r_B}^{r_A} \frac{Q}{4\pi\epsilon_0 r^2} a_r \cdot dr a_r \\ &= \frac{Q}{4\pi\epsilon_0} \left[\frac{1}{r_A} - \frac{1}{r_B} \right] \end{aligned}$$

Then,

$$V_{AB} = V_A - V_B$$

Where, V_B and V_A are the potentials at B and A respectively.

Potential difference produced by a line charge

We know,

$$E = E_{\rho} a_{\rho} = \frac{\rho_L}{2\pi\epsilon_0\rho} a_{\rho}$$

The potential difference,

$$\begin{aligned} V_{AB} &= - \int_B^A E \cdot dL \\ &= \int_B^A \left(\frac{\rho_L}{2\pi\epsilon_0\rho} a_{\rho} \right) (d\rho a_{\rho} + \rho d\varphi a_{\varphi} + dz a_z) \\ &= - \frac{\rho_L}{2\pi\epsilon_0} \int_B^A \frac{d\rho}{\rho} \\ &= - \frac{\rho_L}{2\pi\epsilon_0} [\ln \rho]_B^A \\ V_{AB} &= - \frac{\rho_L}{2\pi\epsilon_0} \ln \frac{A}{B} \end{aligned}$$

Work done when displacement (ρ) of source to point charge increase or decrease

We know,

$$E = \frac{\rho_L}{2\pi\epsilon_0\rho} a_{\rho}$$

We also know, work done

$$\begin{aligned} W &= -Q \int_{initial}^{final} E \cdot dL \\ &= -Q \int_{\rho_1}^{\rho_2} \left(\frac{\rho_L}{2\pi\epsilon_0\rho} a_{\rho} \right) \cdot (d\rho a_{\rho} + \rho d\varphi a_{\varphi} + dz a_z) \\ &= - \frac{Q\rho_L}{2\pi\epsilon_0} \int_{\rho_1}^{\rho_2} \frac{d\rho}{\rho} \\ &= - \frac{Q\rho_L}{2\pi\epsilon_0} [\ln \rho]_{\rho_1}^{\rho_2} \\ W &= - \frac{Q\rho_L}{2\pi\epsilon_0} \ln \frac{\rho_2}{\rho_1} \end{aligned}$$

Work done when displacement (ρ) is same but, change in angle (φ).

We know,

$$E = \frac{\rho_L}{2\pi\epsilon_0\rho} a_\rho$$

We also know, work done

$$\begin{aligned} W &= -Q \int_{initial}^{final} E \cdot dL \\ &= -Q \int_{\rho_1}^{\rho_2} \left(\frac{\rho_L}{2\pi\epsilon_0\rho} a_\rho \right) \cdot dL \end{aligned}$$

Here, $dL = d\rho a_\rho + \rho d\varphi a_\varphi + dz a_z$. Change in angle (φ) makes (ρ) and (z) zero.

$$= -Q \int_{\rho_1}^{\rho_2} \left(\frac{\rho_L}{2\pi\epsilon_0\rho} a_\rho \right) \cdot (\rho d\varphi a_\varphi)$$

$$W = 0$$

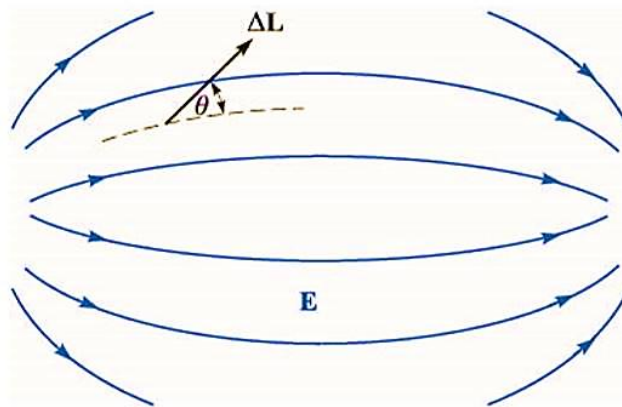
Potential Gradient (W.H.Hayt)

We know

$$V = - \int E \cdot dL$$

For a very short element of length ΔL along which E is constant, leading to an incremental potential difference ΔV ,

$$\Delta V = - E \cdot \Delta L$$



If we designate the angle between ΔL and E as θ , then

$$\Delta V = - E \cdot \Delta L \cos\theta$$

Then,

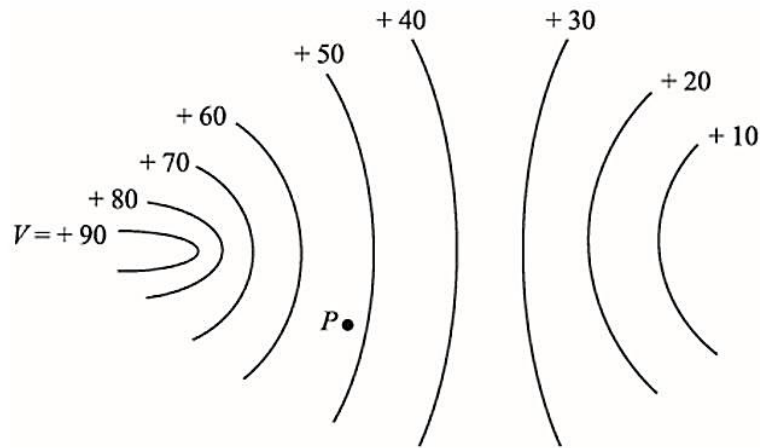
$$\frac{dV}{dL} = - E \cos\theta$$

It is obvious that the maximum positive increment of potential, ΔV_{max} , will occur when $\cos\theta$ is -1 , or ΔL points in the direction opposite to E . For this condition,

$$\left. \frac{dV}{dL} \right|_{max} = E$$

Characteristics of the relationship between E and V :

- The magnitude of the electric field intensity is given by the maximum value of the rate of change of potential with distance.
- This maximum value is obtain when the direction of the distance increment is opposite to E .



*Equipotential surfaces shown as lines in the two dimensional sketch.

At P, small incremental distance ΔL in various directions, to find that direction in which the potential is changing the most rapidly. From the figure this direction appears to be left and slightly upward. So the electric field intensity is therefore oppositely directed (to the right and slightly downward at P). Its magnitude is given by dividing the small increase in potential by the small element of length.

The direction in which the potential is increasing the most rapidly is perpendicular to the equipotentials (in the direction of increasing potential). If ΔL is directed along an equipotential, $\Delta V = 0$. Then,

$$\Delta V = -E \cdot \Delta L = 0$$

Since neither E nor ΔL is zero, E must be perpendicular to ΔL or equipotentials.

Now, by letting a_N be a unit vector normal to the equipotential surface and directed toward the higher potentials. The electric field intensity is then expressed in terms of the potential,

$$E = - \left. \frac{dV}{dL} \right|_{max} a_N$$

The magnitude of E is given by the maximum space rate of change of V and the direction of E is normal to the equipotential surface (in the direction of decreasing potential).

Since $dV/dL|_{max}$ occurs when ΔL is in the direction of a_N ,

$$\left. \frac{dV}{dL} \right|_{max} = \frac{dV}{dN}$$

$$E = - \frac{dV}{dN} a_N$$

The operation on V by which $-E$ is obtained is known as the gradient and the gradient of a scalar field T is defined as

$$\text{Gradient of } T = \text{grad } T = \frac{dT}{dN} a_N$$

Using the new term, we now may write

$$E = -\text{grad } V$$

Now,

$$dV = \frac{\delta V}{\delta x} dx + \frac{\delta V}{\delta y} dy + \frac{\delta V}{\delta z} dz$$

Also,

$$V = - \int E \cdot dL$$

$$dV = -E \cdot dL = -E_x dx - E_y dy - E_z dz$$

Since both expressions are true for any dx, dy and dz , then

$$E_x = - \frac{\delta V}{\delta x}$$

$$E_y = - \frac{\delta V}{\delta y}$$

$$E_z = - \frac{\delta V}{\delta z}$$

Then,

$$E = - \left(\frac{\delta V}{\delta x} a_x + \frac{\delta V}{\delta y} a_y + \frac{\delta V}{\delta z} a_z \right)$$

$$E = - \left(\frac{\delta}{\delta x} a_x + \frac{\delta}{\delta y} a_y + \frac{\delta}{\delta z} a_z \right) \cdot V$$

$$\therefore E = -\nabla \cdot V \quad \text{where, } \left[\nabla = \frac{\delta}{\delta x} a_x + \frac{\delta}{\delta y} a_y + \frac{\delta}{\delta z} a_z \right]$$

$$\nabla V = \frac{\delta V}{\delta x} a_x + \frac{\delta V}{\delta y} a_y + \frac{\delta V}{\delta z} a_z \quad (\text{Cartesian})$$

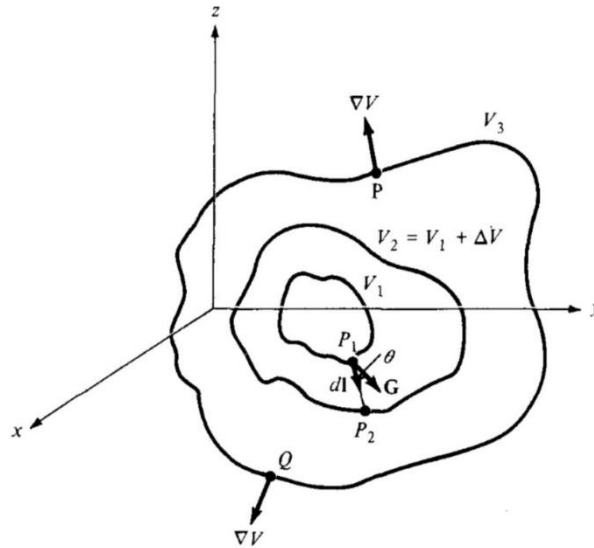
$$\nabla V = \frac{\delta V}{\delta \rho} a_\rho + \frac{1}{\rho} \frac{\delta V}{\delta \varphi} a_\varphi + \frac{\delta V}{\delta z} a_z \quad (\text{Cylindrical})$$

$$\nabla V = \frac{\delta V}{\delta r} a_r + \frac{1}{r} \frac{\delta V}{\delta \theta} a_\theta + \frac{1}{r \sin \theta} \frac{\delta V}{\delta \varphi} a_\varphi \quad (\text{Spherical})$$

Gradient of a Scalar (Sadiku)

The gradient of a scalar field V is a vector that represents both the magnitude and the direction of the maximum space rate of increase of V .

Mathematical expression for the gradient can be obtained by evaluating the difference in the field dV between points P_1 and P_2 of following figure,



$$dv = \frac{\delta V}{\delta x} dx + \frac{\delta V}{\delta y} dy + \frac{\delta V}{\delta z} dz$$

$$= \left(\frac{\delta V}{\delta x} a_x + \frac{\delta V}{\delta y} a_y + \frac{\delta V}{\delta z} a_z \right) \cdot (dx a_x + dy a_y + dz a_z)$$

For convenience, let

$$G = \frac{\delta V}{\delta x} a_x + \frac{\delta V}{\delta y} a_y + \frac{\delta V}{\delta z} a_z$$

Then

$$dV = G \cdot dl$$

$$dV = G \cos\theta dl$$

Or

$$\frac{dV}{dl} = G \cos\theta$$

Where dl is the differential displacement from P_1 to P_2 and θ is the angle between G and dl .

dV/dl is maximum when $\theta = 0$, that is when dl is in the direction of G . Hence,

$$\left. \frac{dV}{dl} \right|_{max} = \frac{dV}{dn} = G$$

Where, dV/dn is the normal derivative. Thus G has its magnitude and direction as those of the maximum rate of change of V. By definition, G is the gradient of V. Therefore,

$$\text{grad } V = \nabla V = \frac{\delta V}{\delta x} a_x + \frac{\delta V}{\delta y} a_y + \frac{\delta V}{\delta z} a_z$$

The gradient of V can be expressed in Cartesian, cylindrical and spherical coordinates.

$$\nabla V = \frac{\delta V}{\delta x} a_x + \frac{\delta V}{\delta y} a_y + \frac{\delta V}{\delta z} a_z \quad (\text{Cartesian})$$

$$\nabla V = \frac{\delta V}{\delta \rho} a_\rho + \frac{1}{\rho} \frac{\delta V}{\delta \varphi} a_\varphi + \frac{\delta V}{\delta z} a_z \quad (\text{Cylindrical})$$

$$\nabla V = \frac{\delta V}{\delta r} a_r + \frac{1}{r} \frac{\delta V}{\delta \theta} a_\theta + \frac{1}{r \sin \theta} \frac{\delta V}{\delta \varphi} a_\varphi \quad (\text{Spherical})$$

Mathematical problem-2:

Given,

$$v = 100 r^2 \sin\theta$$

Electric field, $E = ?$ **Solution:**

We know,

$$\begin{aligned} E &= -\nabla \cdot V \\ \Rightarrow & -\left(\frac{\delta}{\delta r} a_r + \frac{1}{r} \frac{\delta}{\delta \theta} a_\theta + \frac{1}{r \sin\theta} \frac{\delta}{\delta \varphi} a_\varphi\right) \cdot (100 r^2 \sin\theta) \\ \Rightarrow & -\frac{\delta(100 r^2 \sin\theta)}{\delta r} a_r - \frac{\delta(100 r^2 \sin\theta)}{r \delta \theta} a_\theta - 0 \\ \Rightarrow & -(100 \sin\theta) \frac{\delta}{\delta r} (r)^2 a_r - 100r \frac{\delta}{\delta \theta} (\sin\theta) a_\theta \\ \therefore E &= -100 \sin\theta (2r) a_r - 100r \cos\theta a_\theta \quad (\text{Ans}) \end{aligned}$$

Mathematical problem-3:

Given,

$$v = 100 \rho^2$$

Electric field, $E = ?$ **Solution:**

We know,

$$\begin{aligned} E &= -\nabla \cdot V \\ \Rightarrow & -\left(\frac{\delta}{\delta \rho} a_\rho + \frac{1}{\rho} \frac{\delta}{\delta \varphi} a_\varphi + \frac{\delta}{\delta z} a_z\right) \cdot (100 \rho^2) \\ \Rightarrow & -\frac{\delta(100 \rho^2)}{\delta \rho} a_\rho - 0 - 0 \\ \Rightarrow & -100 (2\rho) a_\rho \\ \therefore E &= -200\rho a_\rho \quad (\text{Ans}) \end{aligned}$$

Mathematical problem-4:

Potential field, $V = 2x^2y - 5z$ and a point $P(-4, 3, 6)$. Find following numerical values at point P:

- (a) The potential, V
- (b) The electric field intensity, E
- (c) The direction of E
- (d) The electric flux density D and
- (e) The volume charge density ρ_v

Solution:

- (a) The potential at $P(-4, 3, 6)$

$$V_P = 2(-4)^2(3) - 5(6) = 66 \text{ V}$$

- (b) The electric field intensity

$$E = -\nabla V$$

$$E = -4xy a_x - 2x^2 a_y + 5 a_z \text{ V/m}$$

The value of E at point P is

$$E_P = 48 a_x - 32 a_y + 5 a_z \text{ V/m}$$

And

$$|E_P| = \sqrt{48^2 + (-32)^2 + 5^2}$$

$$|E_P| = 57.9 \text{ V/m}$$

- (c) The direction of E at point P is given by the unit vector

$$a_{E,P} = \frac{48 a_x - 32 a_y + 5 a_z}{57.9}$$

$$a_{E,P} = 0.829 a_x - 0.553 a_y + 0.086 a_z$$

- (d) Assuming these fields exist in free space, then the electric flux density

$$D = \epsilon_0 E = -35.4 xy a_x - 17.71 x^2 a_y + 44.3 a_z \text{ pC/m}^3$$

- (e) The volume charge density

$$\rho_v = \nabla \cdot D$$

$$\rho_v = -35.4 y \text{ pC/m}^3$$

At P,

$$\rho_v = -106.2 \text{ pC/m}^3$$

The Dipole

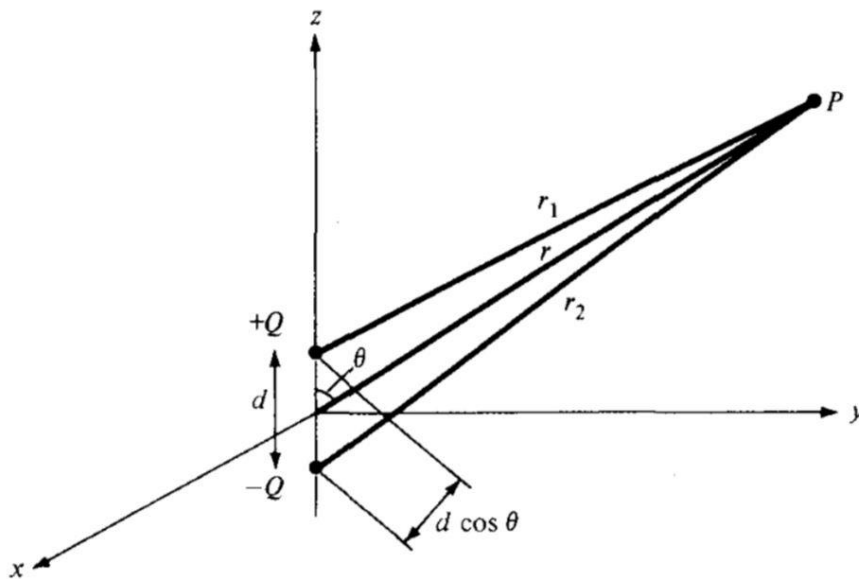
An electric dipole is formed when two point charges of equal magnitude but opposite sign are separated by a small distance.

Consider the dipole shown in following figure, the potential at point $P (r, \theta, \varphi)$ is

$$V = \frac{Q}{4\pi\epsilon_0} \left[\frac{1}{r_1} - \frac{1}{r_2} \right]$$

$$V = \frac{Q}{4\pi\epsilon_0} \left[\frac{r_2 - r_1}{r_1 r_2} \right]$$

Where, r_1 and r_2 are the distances between P and $+Q$ and P and $-Q$ respectively.



If $r \gg d$, $r_2 - r_1 \approx d \cos \theta$ and $r_1 r_2 \approx r^2$ then we have,

$$V = \frac{Q}{4\pi\epsilon_0} \frac{d \cos \theta}{r^2}$$

Electric field intensity,

$$E = -\nabla \cdot V$$

$$E = -\left(\frac{\delta V}{\delta r} a_r + \frac{1}{r} \frac{\delta V}{\delta \theta} a_\theta + \frac{1}{r \sin \theta} \frac{\delta V}{\delta \varphi} a_\varphi \right)$$

Then

$$E = -\left(-\frac{Qd \cos \theta}{4\pi\epsilon_0 r^3} a_r - \frac{Qd \sin \theta}{4\pi\epsilon_0 r^3} a_\theta \right)$$

$$E = -\frac{Qd}{4\pi\epsilon_0} \left(-2 \cos \theta \frac{1}{r^3} a_r - \sin \theta \frac{1}{r^3} a_\theta \right)$$

$$E = -\frac{Qd}{4\pi\epsilon_0 r^3} (-2 \cos \theta a_r - \sin \theta a_\theta)$$

The potential field of the dipole may be simplified by making use of the dipole moment. The vector length directed from $-Q$ to $+Q$ as d and then define the dipole moment as Qd and assign it the symbol P . Thus

$$p = Qd$$

The units of P are Cm .

Since $d \cdot a_r = d \cos\theta$, then we have

$$V = \frac{Q}{4\pi\epsilon_0} \frac{d \cos\theta}{r^2}$$

$$V = \frac{p \cdot a_r}{4\pi\epsilon_0 r^2}$$

This may be generalized as

$$V = \frac{1}{4\pi\epsilon_0} \frac{1}{|r - r'|^2} p \cdot \frac{r - r'}{|r - r'|}$$

$$V = \frac{p \cdot (r - r')}{4\pi\epsilon_0 |r - r'|^3}$$

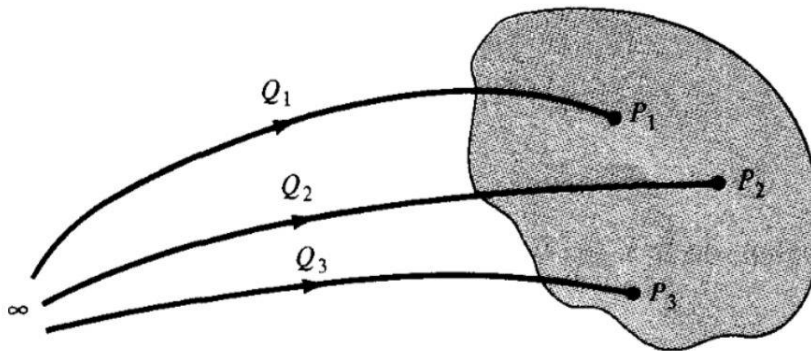
Energy Density in Electrostatic Fields

Three point charges Q_1, Q_2 and Q_3 in an empty space shown in following figure.

- No work is required to transfer Q_1 from infinity to P_1 because the space is initially charge free and there is no electric field.
- The work done in transferring of Q_2 from infinity to P_2 is equal to the product of Q_2 and the potential V_{21} at P_2 due to Q_1 .
- Similarly, the work done in positioning Q_3 at P_3 is equal to $Q_3(V_{32} + V_{31})$, where V_{32} and V_{31} are the potentials at P_3 due to Q_2 and Q_1 respectively.

$$\text{Work to position } Q_2 = Q_2 V_{21}$$

$$\text{Work to position } Q_3 = Q_3 (V_{31} + V_{32})$$



The total work done in positioning the three charges is,

$$W_E = W_1 + W_2 + W_3$$

$$W_E = 0 + Q_2 V_{21} + Q_3 (V_{31} + V_{32}) \quad \text{_____ (1)}$$

If the charges were positioned in reverse order, then,

$$W_E = W_3 + W_2 + W_1$$

$$W_E = 0 + Q_2 V_{23} + Q_1 (V_{12} + V_{13}) \quad \text{_____ (2)}$$

Where, V_{23} is the potential at P_2 due to Q_3 , V_{12} and V_{13} are respectively the potentials at P_1 due to Q_2 and Q_3 .

Now adding equations (1) and (2) gives,

$$2W_E = Q_1 (V_{12} + V_{13}) + Q_2 (V_{21} + V_{23}) + Q_3 (V_{31} + V_{32})$$

$$2W_E = Q_1 V_1 + Q_2 V_2 + Q_3 V_3 \quad [\text{Here, } V_{12} + V_{13} = V_1]$$

$$W_E = \frac{1}{2} Q_1 V_1 + Q_2 V_2 + Q_3 V_3$$

Where, V_1, V_2 and V_3 are total potentials at P_1, P_2 and P_3 respectively. In general, if there are n point charges then,

$$W_E = \frac{1}{2} \sum_{m=1}^{m=n} Q_m V_m \quad (\text{in joules})$$

If, instead of point charges, the region has continuous charge distribution, then we have,

$$W_E = \frac{1}{2} \int \rho_L V dl \quad (\text{Line charge})$$

$$W_E = \frac{1}{2} \int \rho_S V dS \quad (\text{Surface charge})$$

$$W_E = \frac{1}{2} \int \rho_v V dv \quad (\text{Volume charge})$$

Since,

$$\rho_v = \nabla \cdot D$$

For volume charge, we can write

$$W_E = \frac{1}{2} \int_v (\nabla \cdot D) V dv$$

But for any vector A and scalar, the identity

$$\nabla \cdot VA = A \cdot \nabla V + V(\nabla \cdot A)$$

Or,

$$(\nabla \cdot A)V = \nabla \cdot VA - A \cdot \nabla V$$

Applying the identity, we get,

$$W_E = \frac{1}{2} \int_v (\nabla \cdot VD) dv - \frac{1}{2} \int_v (D \cdot \nabla V) dv$$

Applying divergence theorem to the first term on the right-hand side of this equation, we have

$$W_E = \frac{1}{2} \oint_S (VD) \cdot dS - \frac{1}{2} \int_v (D \cdot \nabla V) dv$$

As we know that

- V varies as $1/r$ and D as $1/r^2$ for point charges
- V varies as $1/r^2$ and D as $1/r^3$ for dipoles

So, here in the first term on the right-hand side of the equation

- VD must vary at least as $1/r^3$ and
- dS varies as r^2

Consequently, the first integral of the equation must tend to zero as the surface S becomes large. So the equations reduces to

$$W_E = -\frac{1}{2} \int_v (\mathbf{D} \cdot \nabla V) dv$$

Since, $E = -\nabla V$ and $D = \epsilon_0 E$, we have,

$$W_E = \frac{1}{2} \int \mathbf{D} \cdot \mathbf{E} dv$$

$$W_E = \frac{1}{2} \int (\epsilon_0 \mathbf{E} \cdot \mathbf{E}) dv$$

$$W_E = \frac{1}{2} \int \epsilon_0 E^2 dv$$

From this, we can define electrostatic energy density W_E (in J/m^3) as

$$w_E = \frac{dW_E}{dv} = \frac{1}{2} \mathbf{D} \cdot \mathbf{E} = \frac{1}{2} \epsilon_0 E^2 = \frac{D^2}{2\epsilon_0}$$